

# Schrödinger & Heisenberg Pictures / Ehrenfest th<sup>n</sup>

→  $\hat{H} \rightarrow$  Time independent,

$$U(t, t_0) = e^{\frac{-i}{\hbar} H(t-t_0)}$$

Heisenberg Operators

$$\hat{A}_H = U^\dagger(t, 0) \hat{A}_S U(t, 0)$$

1) At  $t=0$ ,  $\hat{A}_H(0) = \hat{A}_S$

$$\hat{I}_H = U^\dagger(t, 0) \hat{I}_S U(t, 0) = \hat{I}_S$$

2)  $\hat{C}_S = \hat{A}_S \hat{B}_S \rightarrow \hat{C}_H(t) = \hat{A}_H(t) \hat{B}_H(t)$

3)  $[\hat{A}_S, \hat{B}_S] = \hat{C}_S \rightarrow [\hat{A}_H(t), \hat{B}_H(t)] = \hat{C}_H(t)$

4)  $\langle \hat{A}_S \rangle = \langle \hat{A}_H(t) \rangle$

$$\Rightarrow \boxed{\langle \psi, t | \hat{A}_S | \psi, t \rangle = \langle \psi, 0 | \hat{A}_H(t) | \psi, 0 \rangle}$$

$$i\hbar \frac{d\hat{A}_H(t)}{dt} = [\hat{A}_H(t), \hat{H}_H(t)] + i\hbar \left( \frac{\partial \hat{A}_S(t)}{\partial t} \right)_H$$

$\hat{A}_S$  without  
time dep  
time  $\rightarrow$

$$i\hbar \frac{d\hat{A}_H(t)}{dt} = [\hat{A}_H(t), \hat{H}_H(t)]$$

$$\Rightarrow i\hbar \frac{d\langle \hat{A}_H(t) \rangle}{dt} = \langle [\hat{A}_H(t), \hat{H}_H(t)] \rangle$$

$$i\hbar \frac{d\langle \hat{A}_S(t) \rangle}{dt} = \langle [\hat{A}_S(t), \hat{H}_S(t)] \rangle$$

# Example 1 Harmonic Oscillator.

$$\frac{d}{dt} A_H(t) = \frac{1}{i\hbar} [A_H(t), H_H(t)]$$

$$H_S = \frac{\hat{p}^2}{2m} + \frac{1}{2} m\omega^2 \hat{x}^2$$

↓

$$H_H = \frac{\hat{p}_H(t)^2}{2m} + \frac{1}{2} m\omega^2 \hat{x}_H(t)^2$$

if  $H_S(\hat{p}, \hat{x}; t)$   
then  
 $H_H(\hat{p}_H(t), \hat{x}_H(t); t)$

→ Consider  $\hat{x}_S$  &  $\hat{p}_S$  & use them in Heisenberg eq<sup>n</sup> of motion

$$\frac{d}{dt} \hat{A}_H(t) = \frac{1}{i\hbar} [\hat{A}_H(t), \hat{H}_H(t)]$$

for  $\hat{x}_H(t)$

$$\begin{aligned} \frac{d}{dt} \hat{x}_H(t) &= \frac{1}{i\hbar} \left[ \hat{x}_H(t), \frac{\hat{p}_H(t)^2}{2m} + \frac{1}{2} m\hat{x}_H(t)^2 \omega^2 \right] \\ &= \frac{1}{i\hbar} \left[ \hat{x}_H(t), \frac{\hat{p}_H(t)^2}{2m} \right] + \underbrace{\left[ \hat{x}_H(t), \frac{1}{2} m\hat{x}_H(t)^2 \omega^2 \right]}_0 \end{aligned}$$

$$= \frac{1}{2m i\hbar} [A, B C] = [A, B] C + B [A, C]$$

$$= \frac{1}{2m i\hbar} \left[ \hat{x}_H(t), \hat{p}_H(t) \hat{p}_H(t) \right]$$

$$= \frac{1}{2m i\hbar} \left[ \underbrace{[\hat{x}_H(t), \hat{p}_H(t)]}_{i\hbar} \hat{p}_H(t) + \hat{p}_H(t) [\hat{x}_H(t), \hat{p}_H(t)] \right]$$

$$= \frac{2i\hbar}{2m i\hbar} \hat{p}_H(t) = \frac{\hat{p}_H(t)}{m}$$

$$\Rightarrow \frac{d}{dt} \hat{x}_H(t) = \frac{\hat{p}_H(t)}{m}$$

⊙  
→  $E_H \text{ for } H_H$

$$m \frac{d}{dt} \langle \hat{x} \rangle = \langle \hat{p}_x \rangle$$



$$\therefore \hat{x}_H = \hat{x}_S \cos \omega t + \frac{1}{m\omega} \hat{p} \sin \omega t$$

$$\hat{p} =$$

$$\hat{p}_H = \hat{p} \cos \omega t - m\omega \hat{x} \sin \omega t.$$

## Creation & Annihilation operators.

→ What are the Heisenberg operators corresponding to simple harmonic oscillator creation & annihilation operator.

$$\hat{a}_S \Rightarrow \hat{a}_H(t)$$

$$\hat{a}_S^\dagger \Rightarrow \hat{a}_H^\dagger(t)$$

S.H.O Hamiltonian → Time independent

$$\begin{aligned} \Rightarrow \text{We can use } \hat{A}_H(t) &= U^\dagger(t,0) \hat{A}_S U(t,0) \\ &= e^{+\frac{i}{\hbar} H t} \hat{A}_S e^{-\frac{i}{\hbar} H t} \end{aligned}$$

$$\hat{H} = \hbar \omega \left( \underbrace{\hat{a}^\dagger \hat{a}}_{\hat{N}} + \frac{1}{2} \right)$$

Additive constant has no effect on the commutator.

How to evaluate  $\hat{a}_H(t)$ ?

$$\hat{a}_H(t) = e^{+\frac{i}{\hbar} \hat{N} t} \hat{a}_S e^{-\frac{i}{\hbar} \hat{N} t} = e^{i \omega t} \hat{a} \quad H = \underbrace{(\hat{a}^\dagger \hat{a})}_{\hat{N}} \hbar \omega$$

$$\begin{aligned} \frac{d}{dt} \hat{a}_H(t) &= \omega \hat{a}_S e^{i \omega t} [\hat{N}, \hat{a}] e^{-i \omega t} \\ &= \omega e^{i \omega t} [\hat{N}, \hat{a}] e^{-i \omega t} \\ &= -\omega e^{i \omega t} \hat{a}_S e^{-i \omega t} \end{aligned}$$

$$\therefore \frac{d}{dt} \hat{a}_H(t) = -\hat{a}_H(t) i \omega$$

$$\hat{a}(t=0) = \hat{a}$$

$$\text{Sol}^n : \hat{a}_H(t) = \hat{a} e^{-i \omega t}$$

$$\hat{a}_H^\dagger(t) = \hat{a}^\dagger e^{i \omega t}$$