Rohan Kulkarni

June 28, 2021

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Why Long-Lived-Particles (LLPs)?

Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

LLP specific detectors FASER MATHUSLA BELLE II

1 Why Long-Lived-Particles (LLPs)?

2 Lifetime of a particle

3 Signature of "Long-lived particles" at Collider experiments

Direct detection Indirect detection

4 LLP specific detectors FASER MATHUSLA BELLE II

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Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

LLP specific detectors FASER MATHUSLA BELLE II

Conclusion

Why Long-Lived-Particles (LLPs)?

Motivation for LLP

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Conclusion

LLPs?

• LLPs ($\tau \sim ns, c\tau \sim cm$) : essential part of SM \Rightarrow Reason to believe for them to be in BSM

Dark-matter as LLPs?

- Models of WIMP DM → Null results to date in indirect detection (ID), direct detection (DD), and missing energy searches
 - WIMP DM \rightarrow Severely constrained regions of parameter space
 - Broader investigation into possible signals of particle dark matter

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Lifetime of a particle

Decay rates

- Large collection : *decaying particles* at time t : N(t).
- Decay rate Γ : probability per unit time \rightarrow given particle disintegrate
- Rate : particles decrease

$$\frac{dN}{dt} = -\Gamma N \Rightarrow N(t) = N(0) e^{-\Gamma t}$$

• Mean lifetime \rightarrow reciprocal of the decay rate

$$\tau = \frac{1}{\Gamma}$$

- Reality : most particles \rightarrow decay by several different routes.
 - Total decay rate \rightarrow sum of individual rates and so is their lifetimes

$$\Gamma_{\mathrm{tot}} = \sum_{k=1}^{n} \Gamma_{k} \Rightarrow \tau = rac{1}{\Gamma_{\mathrm{tot}}}$$

• Branching ratios for k'th decay mode : $\frac{\Gamma_k}{\Gamma_{tot}}$

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Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Multiple DM theories predicting LLPs

		Small coupling	Small phase space	Scale suppression
SUSY	GMSB			\checkmark
	AMSB		\checkmark	
	Split-SUSY			✓
	RPV	\checkmark		
NN	Twin Higgs	\checkmark		
	Quirky Little Higgs	\checkmark		
	Folded SUSY		\checkmark	
DM	Freeze-in	\checkmark		
	Asymmetric			 ✓
	Co-annihilation		\checkmark	
Portals	Singlet Scalars	\checkmark		
	ALPs			\checkmark
	Dark Photons	\checkmark		
	Heavy Neutrinos			\checkmark

Figure: Table of theories predicting LLPs[1]

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Lifetime of a particle

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DM particle SM particle



Figure: Toy diagram of a freeze-in scenario

- Feeble coupling constant $y_{\chi} \rightarrow Making \chi$ thermally decoupled from the plasma
- This *feebleness* \Rightarrow **long lifetime** of A

Freeze-in DM (FIMPs)

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Freeze-in DM : Decay rate

• **Relic abundance** of χ is related to the A decay width Γ_A by



 g_* is the number of relativistic degrees of freedom at temperatures $\mathcal{T} \approx m_2$ (around the A mass)

- In the SM $g_{*}~(100$ GeV) $\simeq 100$ and $g_{*}~(100$ MeV) $\simeq 10$
- Assuming $\chi
 ightarrow$ all DM today i.e. $\Omega_\chi h^2 = 0.11
 ightarrow$ inverse decay width of A as

$$\Gamma^{-1}\left(A
ightarrow \chi + B_{\mathsf{SM}}
ight) \sim \left(rac{m_1}{100 \mathrm{GeV}}
ight) \left(rac{200 \mathrm{GeV}}{m_2}
ight)^2 \left(rac{100}{g_*\left(m_2
ight)}
ight)^{3/2} imes 10^6 \mathrm{~ns} \sim 0.01 \mathrm{~secs}$$

Co-annihiliating DM

- **DM relic abundance** → annihilation between two different species
- f : SM particle, ψ : BSM LLP, χ : DM particle
- Long lifetime \rightarrow set by a suppressed phase space



Figure: Feynman diagram for a co-annihilation scenario[1]

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Asymmetric DM

- DM particle \neq own antiparticle : relic dark matter density \rightarrow particle-antiparticle asymmetry (Like the Baryon asymmetry)
 - DM production \rightarrow non-thermally (out-of-equilibrium process)
- Easiest scenario for DM production using asymmetry
 - Early universe : Particle species $\psi o m_\psi > m_\chi$ with an abundance $\Omega_\psi o$ decays to χ

$$\Omega_\chi\simeq\Omega_\psirac{m_\chi}{m_\psi}$$



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Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Conclusion

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Lifetime of a particle

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Detecting Dark-matter



Figure: [3]

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Outgoing

Particle

Figure: Not to scale transverse schematic of a typical collider [1]

Kinematics in a detector

Inner Tracking Detector (ID)

Electromagnetic Calorimeter (ECAL)

Hadronic Calorimeter (HCAL)

Muon System (MS)

Kinematics of LLPs [1]

•
$$d_T = \beta \gamma c \tau$$
, $\gamma = \frac{E}{m} = \frac{1}{\sqrt{1-\beta^2}}$,
 $\beta = \frac{v}{c} = \frac{|\vec{p}_T|}{E}$
• E : Calorimeter, $|\vec{p}_T|$: Track bence

•
$$N(t) = N_0 e^{-\frac{t}{\tau}}$$

•
$$P_{\text{dec}} = \frac{1}{4\pi} \int_{\Delta\Omega} d\Omega \int_{L_1}^{L_2} \frac{1}{d} e^{-\frac{L}{d}}$$

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Figure: Cross section of a collider[4]

- Particles produce ightarrow region of ionization ightarrow solid-state / gaseous detectors ightarrow hits.
 - Fit into trajectory \rightarrow track \rightarrow Get charge / momentum

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Decays within the tracker



Figure: Fraction of LLPs decaying \rightarrow Left : within 30 cm, Right : between 30 - 100 cm [2]

- Most theories : Predict lifetimes of LLP >> 25 ns.
- Detect them in conventional subsystems $^{\sf due\ to} \longrightarrow {\sf Exponential}$ decay probability

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Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Direct detection

Indirect detection

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Direct detection : Anomalous Ionization

- CLLP \rightarrow Leaves *track* in ID \Rightarrow Direct detection possible
 - If $m_{\text{CLLP}} > m_{\text{proton}} \rightarrow \text{Produced with lower } \beta$ (Compared to : Track forming SM particle)
 - Detect : Slow moving / heavily charged particle \rightarrow Anomalously large $\langle \frac{dE}{dx} \rangle$
 - Bethe-Bloch formula $\left\langle \frac{dE}{dx} \right\rangle \sim -\frac{z^2}{\beta^2} \cdot \left[\ln \left(\frac{\beta^2}{(1-\beta^2)} \right) \beta^2 + C \right]$ (Ionization energy lost per unit distance traveled)



Figure: Anomalous ionization of a heavy CLLP[1]

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Lifetime of a particle

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Indirect detection I : Displaced Tracks



Figure: Displaced track vs Prompt track[1]

- $\bullet~\mbox{Neutral LLP} \rightarrow \mbox{Transverses}$ some macroscopic distance within ID
 - Decays into charged particle/s →Leaves a *displaced track* or *a displaced vertex* (next slide)

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Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Indirect detection II : Displaced Vertices



Figure: Displaced vertex vs Prompt vertex[1]

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Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

LLP specific detectors FASER MATHUSLA BELLE II

Conclusion

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Why Long-Lived-Particles (LLPs)?

Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

LLP specific detectors

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Conclusion

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Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

LLP specific detectors

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Figure: Left : Location of FASER w.r.t Atlas, Right : View of FASER in a tunnel[7]

- Goal : Detect LLPs / decay products ightarrow Transversed $\sim 150m$ (Inagurated May 2021)
- Isolation \rightarrow Low SM background (Most SM background : near the ATLAS IP)
- New resolution/parameter space \rightarrow LLP detection.

MATHUSLA



Figure: Schematic of proposed MATHUSLA detector[9]

- $200 \times 200 \times 20 \text{ m}^3$ in size, roughly 100 m above CMS/ATLAS caverns.
- Neutral LLPs : very large lifetimes produced in the collisions \to decay within the volume of MATHUSLA \to displaced vertices could be reconstructed

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Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Belle II (e^+e^- collider)



Figure: Schematic of Belle II electron-positron collider[8]

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Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

LLP specific detectors FASER MATHUSLA BELLE II

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Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Displaced vertex signatures



Figure: [10]

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Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Figure: [10]

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- LLPs \rightarrow Natural prediction by many theories.
 - LLP searches important : It is a *strong DM candidate*, also to cover entire *spectrum of DM candidates*
- Increase chances of LLP detection,
 - Low SM backgrounds
 - Extra detectors : *far-distance* from collision point \rightarrow *decay products* get an opportunity to be detected
- Very young field : huge potential for discovering different aspects of BSM, both theoretically and experimentally

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Why Long-Lived-Particle (LLPs)?

Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

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Conclusion

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Why Long-Lived-Particles (LLPs)?

Lifetime of a particle

Signature of "Long-lived particles" at Collider experiments Direct detection Indirect detection

LLP specific detectors FASER MATHUSLA BELLE II

Conclusion

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