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**Problem.  $HZ$  production at a  $e^+e^-$  collider**

Consider the process

$$e^-(p_1)e^+(p_2) \rightarrow h(k_1)Z(k_2) \quad (1)$$

with  $q = p_1 + p_2$ . The Feynman rule for  $Zf\bar{f}$ -vertex is

$$-i \frac{g}{4 \cos \theta_W} \gamma_\mu (V - A\gamma_5), \quad V = 2I_W^3 - 4Q \sin^2 \theta_W, \quad A = 2I_W^3. \quad (2)$$

being  $I_W^3 = \pm 1/2$  the isospin of the fermion and  $Q$  its electric charge. In the electron case ( $f = e^-$ ):

$$V = -1 + 2 \sin^2 \theta_W, \quad A = -1. \quad (3)$$

The Feynman rule for the  $hZZ$  vertex is

$$ig \frac{m_Z}{\cos \theta_W} g_{\mu\nu} = i \frac{2m_Z^2}{v} g_{\mu\nu}. \quad (4)$$

- (a) Draw the Feynman diagram for the process and write down the amplitude of the process using the couplings mentioned above. Use unitary gauge.
- (b) Write down the amplitude in terms of  $G_F$  by using

$$\frac{m_W}{m_Z} = \cos \theta_W, \quad \frac{g^2}{8m_W^2} = \frac{G_F}{\sqrt{2}} \quad (5)$$

- (c) Compute the unpolarised squared matrix-element  $|\bar{M}|^2$ , summed over final state spin and averaged over the initial polarizations, taking  $m_e \rightarrow 0$
  - (d) Compute the differential cross section of the process in the C.O.M frame.
  - (e) Compute the total cross section of the above process. Check how the total cross section behaves with variation of  $m_h$ .
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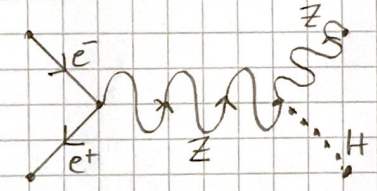
Consider :  $e^-(p_1)e^+(p_2) \rightarrow h(k_1)Z(k_2)$  with  $q = p_1 + p_2$ .

Zff vertex :  $-i \frac{g}{4 \cos \theta_w} \gamma_\mu (V - A \gamma_5)$  ,  $V = 2Z_W^2 - 4Q \sin^2 \theta_w$  ↗ Charge  
 $A = 2I_W^3$  ↘  $I_W = \pm \frac{1}{2}$   
isospin of f.

For (f=e):  $V = -1 + 2 \sin^2 \theta_w$  ,  $A = -1$

hZZ vertex :  $ig \frac{m_Z}{\cos \theta_w} g_{\mu\nu} = i \frac{2m_Z^2}{V} g_{\mu\nu}$

(a) Draw a feynman diagram for the process & Write down amplitude using couplings abv. Use Unitary gauge.

$iM =$    $= \bar{V}(p_2) \frac{-ig}{4 \cos \theta_w} \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} (g_{\mu\nu} - \frac{k^\mu k^\nu}{m_Z^2}) \frac{ig m_Z}{\cos \theta_w} g_{\nu\lambda} \epsilon^\lambda(k_2)$

(b) Write amplitude in terms of  $G_F$  using :  $\frac{m_W}{m_Z} = \cos \theta_w$  ;  $\frac{g^2}{8m_W^2} = \frac{G_F}{\sqrt{2}}$

$\rightarrow m_W = \sqrt{\frac{2}{8G_F}} g \Rightarrow \frac{1}{\cos \theta_w} = \sqrt{\frac{8G_F}{2}} \frac{m_Z}{g}$

$\therefore iM = \bar{V}(p_2) \frac{1}{g} \frac{8G_F}{\sqrt{2}} m_Z^3 \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} (-g^{\mu\nu} + \frac{k^\mu k^\nu}{m_Z^2}) g_{\nu\lambda} \epsilon^\lambda(k_2)$   
 $= \bar{V}(p_2) \sqrt{2} G_F m_Z^3 \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} (-g^{\mu\nu} + \frac{k^\mu k^\nu}{m_Z^2}) g_{\nu\lambda} \epsilon^\lambda(k_2)$   
 $= \bar{V}(p_2) \sqrt{2} G_F m_Z^3 \gamma_\mu (V - A \gamma_5) u(p_1) \frac{-i}{k^2 - m_Z^2} (\frac{k^2}{m_Z^2} - 4) \epsilon^\lambda(k_2)$

(c) Compute  $|M|^2$  : Summed over final spin state & Avg over init. polariz.

$m_e \rightarrow 0$

$|M|^2 = \frac{G_F^2}{2(s - m_Z^2)^2} m_Z^6 \sum_{s,s'} [\bar{V}_s \gamma_\mu (V - A \gamma_5) u_{s'}] (\sum_p \epsilon_p^\mu \epsilon_p^\nu) (\frac{k^2}{m_Z^2} - 4) [\bar{u}_{s'} \gamma_\nu (V - A \gamma_5) V_s]$



$$= \frac{G_F^2}{2(s-m_Z^2)^2} m_Z^6 \left( \frac{k^2}{m_Z^2} - 4 \right) \sum_{s_1 s_1'} \left\{ \left[ \bar{u}_{s_1} \gamma_\mu (V-A\gamma_5) u_{s_1'} \right] \right. \\ \left. \left( \frac{k_2^\mu k_2^\nu}{m_Z^2} - g^{\mu\nu} \right) \left[ u_{s_1'} \gamma_\nu (V-A\gamma_5) v_{s_2} \right] \right\}$$

↓ Cauchy's trick

$$= \frac{G_F^2}{2(s-m_Z^2)^2} m_Z^6 \left( \frac{k^2}{m_Z^2} - 4 \right) \times \left[ \text{Tr} \left[ \not{p}_2 \gamma_\mu (V-A\gamma_5) \not{p}_1 \gamma_\nu (V-A\gamma_5) \right] \left( \frac{k_2^\mu k_2^\nu}{m_Z^2} - g^{\mu\nu} \right) \right]$$

$$\text{Tr} \left[ \not{p}_2 \gamma_\mu (V-A\gamma_5) \not{p}_1 \gamma_\nu \right] = p_2^\lambda p_1^\sigma \text{Tr} \left[ \gamma_\lambda \gamma_\mu (V-A\gamma_5) \gamma_\sigma \gamma_\nu (V-A\gamma_5) \right]$$

$$= p_2^\lambda p_1^\sigma \left\{ V^2 \text{Tr} \left[ \gamma_\lambda \gamma_\mu \gamma_\sigma \gamma_\nu \right] - VA \text{Tr} \left[ \gamma_\lambda \gamma_\mu \gamma_\sigma \gamma_\nu \gamma_5 \right] \right. \\ \left. - VA \text{Tr} \left[ \gamma_\lambda \gamma_\mu \gamma_5 \gamma_\sigma \gamma_\nu \right] + A^2 \text{Tr} \left[ \gamma_\lambda \gamma_\mu \gamma_5 \gamma_\sigma \gamma_\nu \gamma_5 \right] \right\}$$

$$= p_2^\lambda p_1^\sigma \left\{ (V^2 + A^2) (g_{\lambda\mu} g_{\sigma\nu} - g_{\lambda\sigma} g_{\mu\nu} + g_{\lambda\nu} g_{\mu\sigma}) \right. \\ \left. - 8iVA \epsilon_{\lambda\mu\sigma\nu} \right\}$$

$$\Rightarrow g^{\mu\nu} \text{Tr} \left[ \not{p}_2 \gamma_\mu (V-A\gamma_5) \not{p}_1 \gamma_\nu \right]$$

$$= \left\{ (V^2 + A^2) \left[ (p_2 \cdot p_1) - 4(p_2 \cdot p_1) + (p_2 \cdot p_1) \right] \right\}$$

$$= -2(V^2 + A^2) (p_2 \cdot p_1)$$

$$k^\mu k^\nu \text{Tr} \left[ \not{p}_2 \gamma_\mu (V-A\gamma_5) \not{p}_1 \gamma_\nu \right] = \left\{ (V^2 + A^2) \left[ 2(p_1 k_2)(p_2 k_2) - (p_1 p_2) k_2^2 \right] \right\}$$

$$\Rightarrow |M|^2 = \frac{G_F^2}{2(s-m_Z^2)^2} m_Z^6 \left( \frac{k^2}{m_Z^2} - 4 \right) (V^2 + A^2) \left[ \frac{2(p_1 k_2)(p_2 k_2)}{m_Z^2} - (p_1 p_2) \right. \\ \left. + 3(p_1 p_2) \right]$$

$$= \frac{G_F^2}{2(s-m_Z^2)^2} m_Z^6 \left[ \frac{k^2}{m_Z^2} - 4 \right] 2(V^2 + A^2) \left\{ (p_1 p_2) + \frac{(p_1 k_2)(p_2 k_2)}{m_Z^2} \right\}$$

(d) Cross Section in COM frame.

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} \frac{1}{P_i} \times \frac{G_F^2}{2(s-m_Z^2)^2} m_Z^6 \left( \frac{k^2}{m_Z^2} - 4 \right) 2(V^2 + A^2) \left\{ (p_1 p_2) + \frac{(p_1 k_2)(p_2 k_2)}{m_Z^2} \right\}$$



### (3) Width of the W boson.

(a) Write all the allowed 2-body decays of the  $W^-$  boson in the SM. Neglect quark mixing ( $V_{CKM}=1$ )

→ Leptons :

$$W^- \rightarrow e^- \bar{\nu}_e$$

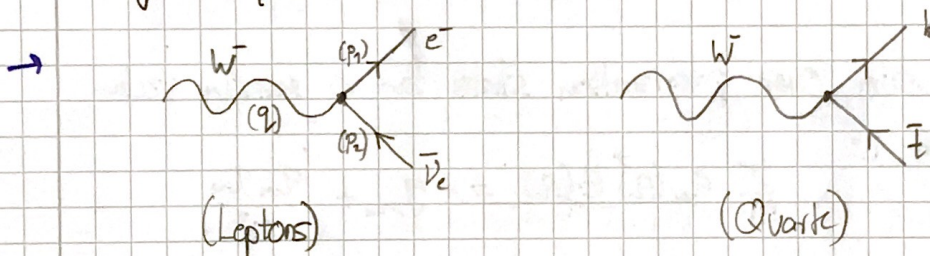
$$W^- \rightarrow \mu^- \bar{\nu}_\mu$$

$$W^- \rightarrow \tau^- \bar{\nu}_\tau$$

Quarks :

$W^- \rightarrow \bar{u}s$	$W^- \rightarrow \bar{u}b$	$W^- \rightarrow \bar{c}s$
$W^- \rightarrow \bar{u}d$	$W^- \rightarrow \bar{c}d$	$W^- \rightarrow \bar{c}b$
$W^- \rightarrow \bar{t}d$	$W^- \rightarrow \bar{t}s$	$W^- \rightarrow \bar{t}b$

(b) Draw the Feynman diagrams for a decay to Lepton & for a decay to quarks.



(c) Neglecting quark mixing, the Feynman rule for  $W\bar{f}f'$  interaction is the same for quarks & leptons :

$$i \frac{g}{\sqrt{2}} \gamma_\mu P_L = i \frac{g}{2\sqrt{2}} \gamma_\mu (1 - \gamma_5)$$

Write down the amplitude for  $W^-(q) \rightarrow e^-(p_1) \bar{\nu}_e(p_2)$

→ Rules :

- Incoming W-boson :  $\epsilon_\mu(q)$
- Electron :  $\bar{u}(p_1)$
- Anti-neutrino :  $v(p_2)$
- Vertex factor :  $-i \frac{g}{\sqrt{2}} \gamma_\mu P_L = i \frac{g}{2\sqrt{2}} \gamma_\mu (1 - \gamma_5)$

$$-iM = \epsilon_\mu(q) \bar{u}(p_1) \left( -i \frac{g}{\sqrt{2}} \gamma^\mu (1 - \gamma_5) \right) v(p_2)$$

$$M = \frac{g}{\sqrt{2}} \epsilon_\mu(q) \bar{u}(p_1) \frac{\gamma^\mu (1 - \gamma_5)}{2} v(p_2)$$



If we take weak-charged current as:

$$j^\mu = \bar{u}(p_1) \gamma^\mu \frac{1}{2} (1 - \gamma_5) v(p_2)$$

$$M = \frac{g}{\sqrt{2}} \epsilon_\mu(q) j^\mu$$

(d) Take the Fermion masses to be zero, & compute the squared matrix element  $|M|^2$  for  $W^- \rightarrow e^- \bar{\nu}_e$  by

→ Averaging over the polarizations of the initial state

→ Summing over spin of final state particles.

Write it in terms of  $G_F$  &  $m_W$   $\left\{ \begin{array}{l} G_F = \frac{g^2}{8m_W^2} \\ \frac{m_W}{m_Z} = \cos\theta_W \end{array} \right.$

& work out the dependence on the particles' momenta, working in  $W$  rest frame.

Hint: The sum over polarization states for a massive vector boson is

$$\sum_{\text{pol}} \epsilon_\mu(q) \epsilon_\nu^*(q) = -g_{\mu\nu} + \frac{q_\mu q_\nu}{m^2}$$

$$\text{Also, } \gamma_\mu \gamma^\nu \gamma^\mu = -2\gamma^\nu$$

→ Start with Lepton current:  $j^\mu = \bar{u}(p_1) \gamma^\mu \frac{1}{2} (1 - \gamma_5) v(p_2)$

• Now use the rest frame of the  $W^-$  boson & also  $m_W \gg m_e$  to neglect  $m_{\text{lepton}}$ .

$$q = (m_W, \vec{0})$$

$$p_1 = ($$

$$|M|^2 = MM^* = \sum \frac{g^2}{2} \epsilon_\mu^{(i)}(q) j^\mu (j^\nu)^* \epsilon_\nu^{(j)*}(q)$$

$$= \frac{g^2}{2} \sum_{i,j} \epsilon_\mu^{(i)}(q) \epsilon_\nu^{(j)*}(q) \sum_{\text{spin}} j^\mu (j^\nu)^*$$



$$\rightarrow |M|^2 = MM^\dagger = \frac{1}{3} \sum_{\text{pol}}^{\text{Avg}} \sum_{\text{final}}^{\text{Sum}} \frac{g^2}{2} \epsilon_\mu^{(a)}(q) \epsilon_\nu^{(a)}(q)^* \int_{(a)}^{\mu} \int_{(a)}^{\nu}$$

$$= \frac{g^2}{6} \left[ -g_{\mu\alpha} + \frac{q_\mu q_\alpha}{m_W^2} \right] \sum_{\sigma} \bar{u}_\sigma(p_1) \gamma^\mu \frac{1}{2} (1-\gamma_5) \underbrace{v_\sigma(p_2) \bar{v}_\sigma(p_2)}_{\sigma \text{ sum} \Rightarrow \text{Dirac sum}} \frac{1}{2} (1-\gamma_5) u_\sigma(p_2)$$

$$= \frac{g^2}{6} \left[ -g_{\mu\alpha} + \frac{q_\mu q_\alpha}{m_W^2} \right] \sum_{\sigma} \bar{u}_\sigma(p_1) \gamma^\mu \frac{1}{2} (1-\gamma_5) \underbrace{(p_2 + m_{e^-})}_{\sigma \text{ sum} \Rightarrow \text{Dirac sum}} \frac{1}{2} (1-\gamma_5) u_\sigma(p_2)$$

$$= \frac{g^2}{6 \cdot 4} \left[ -g_{\mu\alpha} + \frac{q_\mu q_\alpha}{m_W^2} \right] \gamma^\mu (1-\gamma_5) (p_1 + m_{e^-}) \not{p}_2 \gamma^\alpha (1-\gamma_5)$$

$$= \frac{g^2}{24} \left[ -g_{\mu\alpha} + \frac{q_\mu q_\alpha}{m_W^2} \right] \gamma^\mu (1-\gamma_5) \gamma^\beta (p_1)_\beta \gamma^\nu (p_2)_\nu \gamma^\alpha (1-\gamma_5)$$

$$= \frac{g^2}{24} \left[ -g_{\mu\alpha} + \frac{q_\mu q_\alpha}{m_W^2} \right] (p_1)_\beta (p_2)_\nu \text{Tr}[\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\alpha] (1-\gamma_5)$$

$$= \frac{g^2}{24} \text{Tr}[\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\alpha - \gamma^\mu \gamma^\beta \gamma^\nu \gamma^\alpha \gamma_5]$$

$$= \frac{g^2}{24} \left( \text{Tr}[\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\alpha] - \text{Tr}[\gamma^\mu \gamma^\beta \gamma^\nu \gamma^\alpha \gamma_5] \right)$$

$$= \frac{g^2}{24} \left[ 4(g^{\mu\beta} g^{\nu\alpha} - g^{\mu\nu} g^{\beta\alpha} + g^{\mu\alpha} g^{\beta\nu}) \right]$$

$$- 4i \epsilon^{\mu\nu\alpha\beta}$$

↪ Do not contribute.

$$= \frac{g^2}{4} \left[ -g_{\mu\alpha} + \frac{q_\mu q_\alpha}{m_W^2} \right] (p_1)_\beta (p_2)_\nu [g^{\mu\beta} g^{\nu\alpha} - g^{\mu\nu} g^{\beta\alpha} + g^{\mu\alpha} g^{\beta\nu}]$$

$$= \frac{g^2}{4} \left[ -g_{\mu\alpha} + \frac{q_\mu q_\alpha}{m_W^2} \right] \left( (p_1)^\mu (p_2)^\alpha - (p_1)^\alpha (p_2)^\mu + \cancel{g^{\mu\alpha}} (p_1) \cdot (p_2) \right)$$

$$= \frac{g^2}{4} \left[ \cancel{(p_1 \cdot p_2)} + \cancel{(p_2 \cdot p_1)} - 4(p_1 \cdot p_2) + \frac{1}{m_W^2} \left( \cancel{(p_1 \cdot p_2)} (p_2 \cdot q) - \cancel{(p_1 \cdot q)} (p_2 \cdot q) + (q \cdot q) (p_1 \cdot p_2) \right) \right]$$

$$= \frac{g^2}{4} \left[ -4(p_1 \cdot p_2) + \frac{(q \cdot q)(p_1 \cdot p_2)}{m_W^2} \right]$$

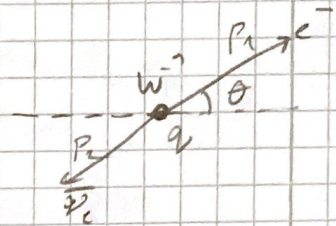


Now take the rest frame for  $W^-$  in  $W^- \rightarrow e^- \bar{\nu}_e$

$$q = (m_W, \vec{0}), \quad p_1 = (E, E \sin \theta, 0, E \cos \theta)$$

$$p_2 = (E, -E \sin \theta, 0, -E \cos \theta)$$

$$\Rightarrow (p_1 \cdot p_2) = 2E^2, \quad (q \cdot q) = m_W^2$$



Gmign

$$|M|^2 = \frac{g^2}{4} \left[ -\cancel{4} \cdot 2E^2 + m_W^2 \frac{2E^2}{m_W^2} \right]$$

$$= -\frac{g^2}{4} \cancel{4} 2E^2 + \frac{g^2}{4} \cancel{m_W^2} \frac{2E^2}{\cancel{m_W^2}}$$

$$= -g^2 2E^2 + g^2 \frac{E^2}{2}$$

$$= -\frac{3}{2} E^2 g^2$$

$$= \left( -\frac{3}{2} \right) g^2 m_W^2$$

Some mistake,  
I should get  
 $\frac{1}{3} g_W^2 m_W^2$

• One can also find that

$$|M_{-1}|^2 = g_W^2 m_W^2 \frac{1}{4} (1 + \cos \theta)^2$$

$$|M_{-1}|^2 = g_W^2 m_W^2 \frac{1}{2} \sin^2 \theta$$

$$|M_{+1}|^2 = g_W^2 m_W^2 \frac{1}{4} (1 - \cos \theta)^2$$

Three possible  $W^-$  polarization modes.

(e) Using this result, ~~compute the total width of~~ write down  $|M|^2$  for  $W^- \rightarrow \bar{u}d$ . What are the changes compared to the leptonic one.

(lowest order)  $\rightarrow$   ~~$\Gamma(W^- \rightarrow e^- \bar{\nu}_e)$~~   
 $\Gamma(W^- \rightarrow \bar{u}d) = 6 \Gamma(W^- \rightarrow e^- \bar{\nu}_e)$

I am assuming due to the heavier weight of quarks & more generations.



(f)



$$d\Gamma = \frac{1}{64\pi^3 m} |\bar{M}|^2 d\Omega$$

$$d\Omega = d\phi d\cos\theta$$

- Compare to scattering formula. ( $2 \rightarrow 2$ )

→ Pretty much the same

$$d\sigma = \frac{1}{64\pi^2 s} |\bar{M}|^2 d\Omega$$

→ We have 2 momenta coming in.

- Compute the partial widths  $W^- \rightarrow e^- \bar{\nu}_e$

$$W^- \rightarrow \bar{u} d$$

$$W^- \rightarrow e^- \bar{\nu}_e$$

$$\int d\Gamma = \int \frac{1}{64\pi^3 m} |\bar{M}|^2 d\Omega = \frac{1}{64\pi^3 m} \int \frac{1}{3} g^2 d\Omega d\Omega$$

$$\Gamma_{e\nu} = \frac{4\pi g^2}{64\pi} = \frac{g^2}{16}$$

$$W^- \rightarrow \bar{u} d$$

$$\Gamma_{ud} = 6 \Gamma_{e\nu} = \frac{3}{8} g^2$$